#### **ORIGINAL PAPER**



# A new semi-analytical approach for self and mutual inductance calculation of hexagonal spiral coil used in wireless power transfer systems

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#### Abstract

Several methods have been proposed in the literature for the calculation of self and mutual inductance. These methods include the use of complex integral analysis, the necessity of having primary and secondary coils with the same dimensions and the limitations of the ratio of the coil dimension to the distance between the coils. To overcome these restrictions, a new semi-analytical estimation method has been proposed in this paper. Calculation the self and mutual inductance by using the same basic formula which is based on Biot Savart Law prevents the formation of complex integrals and helps create a simple solution method. In order to verify the results obtained with the analytical approach, two hexagonal coils with 10 and 20 cm outer side lengths were produced by using litz wire with a conductive cross section of 1.78 mm<sup>2</sup>. The results obtained with the new approach are compared with the finite element analysis, other work presented in the literature and experimental results in order to prove the accuracy of the proposed method.

Keywords Inductance calculation · Wireless power transfer (WPT) · Electric vehicles (EVs) · Inductive power transfer

# 1 Introduction

Interest in wireless power transfer (WPT) systems has been increasing in recent years. Since there is no physical connection between the source and the load in WPT systems, it allows the installation of a simpler and more reliable charging system. The use of WPT systems has already been commercialized in some fields such as electric vehicles (EV) battery charging, mobile phone battery charging, biomedical systems and defense industry applications [1, 2]. Basically, these systems include primary and secondary coils and the

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power is transferred from distance between these coils with the assistance of compensating circuits. In some application ferrite cores are used in order to increase the coupling [3, 4]. In recent years, mostly coreless structures are preferred for light and compact systems and also to prevent core losses. In EV applications the distance between the coils is significantly long which leads to loosely coupled systems. Due to its compact structure, flat spiral coil design is preferred for these types of applications [5–8]. Dimensions and number of turns of coils depend on the power to be transmitted and the operating frequency. In order to determine all system parameters and to obtain the optimum design for WPT systems, it is important to calculate the coil inductance values fast and accurately as they are directly related to the power.

In the literature, there are several proposed methods to calculate the self and mutual inductance, especially for low power applications [9]. The conductor cross sections are of great importance when calculating the inductance values of small sized coils, while it is less important for the large-sized coils [10]. Considering the coil dimensions and air gap for high-power applications with large coil sizes in WPT systems, the copper cross section is not critical in self and mutual inductance calculations. In the inductance calculations of small size spiral coils, the conductor thickness and the distance between the conductors greatly affect the inductance values [11]. However, these smaller effects decrease as the coil dimension increase [12].

In EV applications, it is possible to transfer the desired power from a certain distance with different coil shapes at the appropriate resonance frequency. In addition to conventional coil shapes such as circle, square and rectangle, different coil shapes have also been suggested [7, 8]. In order to calculate the self and mutual inductance, an analytical design of Archimedean spiral coils for EV applications is given and the results are compared with the measurements in [13]. Hexagonal coil shape is one of these different coil types which is generally used as a honeycomb pad in low power applications [14–16]. In high-power applications such as EV's battery charging, hexagonal coil shape is less sensitive to misalignment conditions and is a good candidate for WPT systems [17].

Grover's self and mutual inductance calculation studies are one of the most cited calculation methods in the literature. This method calculates the inductance values quite accurately by using the geometric mean difference (GMD) and arithmetic mean difference (AMD). However, depending on the coil dimensions, the variation of the coefficients in the formulas makes it difficult to use this method. In addition, the formulas are defined only for two identical coils and the short coils which has a small ratio of total coil length to one side length of coil. Therefore, a simple and applicable formula to any coil with any dimensions is needed. Such an equation can be derived by using the Biot-Savart Law to calculate the flux density. Calculation of the self and mutual inductance of multiple coils were carried out with this approach for rectangular coils to obtain an accurate and simple solution, and the results were compared with the measurements in [18]. For the calculation of mutual inductance in hexagonal coil structure, the magnetic flux density generated by a side length in the hexagonal area was calculated [19]. For this purpose, the authors divided the hexagon into 4 triangles and a rectangular area and calculated the amount of flux for each area separately. Then the result is multiplied by six to obtain total amount of flux and calculate the mutual inductance of hexagonal coils. In another study the hexagon is divided into only triangles [20]. For both of these studies, the calculation of the flux amount in the rectangular area is simple, while the calculation of the triangle area is more complex and needs several solution steps. In another proposed method to facilitate the calculation in hexagonal structure, the hexagonal area is divided into small triangles as in the finite element analysis (FEA) and the magnetic flux density in each triangle is considered equal. In this method, the number of meshes should be increased in order to reach the correct result which means more complexity and longer solution time [21]. Greenhouse compared different inductance calculation methods in [22]. There are also monomial formulas derived empirically or large databases in the calculation of self-inductance [23]. In recent studies, multiple coil models are studied to increase magnetic coupling with data transfer capability for WPT systems by using Grover's formula for inductance calculations [24]. A fast and accurate calculation of mutual inductance for dynamic wireless charging method is proposed in [25]. In this work, a novel, fast and basic method for the calculation of self and mutual inductance has been proposed. The method is based on Biot Savart method and has a similar approach to the inductance calculation of the rectangular coils. In order to reduce the complex integrals which are existing due to the triangular areas in the hexagonal inner area, they are also divided into rectangular areas. Thus, only one formula is used throughout the calculation. An algorithm has been developed to find a significant number of division of the triangle areas into rectangular pieces in order to calculate the accurate inductance values. In addition, instead of dividing the entire hexagonal area into small pieces, dividing the triangular areas only into small rectangles will greatly increase the speed of the solution.

In this paper, an inductance estimation approach which can be used for both mutual and self-inductance calculations has been proposed and experimentally proved. In the second section, the relationship between the coil inductances and the transferred power is given by the equations and the self and mutual inductance calculation of a rectangular coil which is the basic approach of the proposed method are explained briefly. In the third section, the equations derived for hexagonal coils are given. The proposed semi-analytical self and mutual inductance calculation method for hexagonal coil is described in detail in section four and compared with other methods in the literature. The comparison of results show that the proposed method gives quite accurate results in both the self and mutual inductance calculations for hexagonal coil structure.

## 2 Fundamental of coil inductance calculations for WPT design

In this section, the importance of inductance calculations for WPT systems and the method proposed for these calculations are explained. Figure 1 shows the loosely coupled WPT system and its equivalent circuit with the Series–Series compensation topology which is the most common topology for EV battery charging applications. According to Eq. (1), the maximum power that can be transferred to the load depends on the resonant frequency ( $\omega$ ), mutual inductance (M), selfinductance ( $L_2$ ) and the quality factor ( $Q_2$ ) of the secondary coil. For the S–S topology the quality factor equation for secondary coil is given in (2).

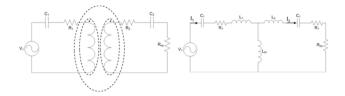


Fig. 1 The loosely coupled WPT system and its equivalent circuit with the Series–Series compensation topology

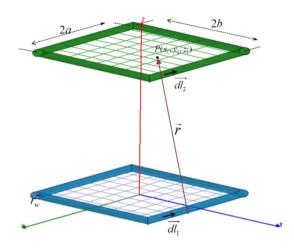


Fig. 2 Single turn rectangular coils

$$P_0 = \frac{\omega M^2 Q_2}{L_2} I_1^2$$
 (1)

$$Q_2 = \frac{\omega L_2}{R_{\rm eq}} \tag{2}$$

A design process to determine the optimum parameters of a WPT system such as operating frequency, number of turns, the wire section with a lowest copper mass is given in [26] for EV applications. Inductance calculations based on Neumann's expression for rectangular coils are also given in the same paper. As claimed in the paper, the calculation of the inductance values is the one of the most important parts of the design. When the self and mutual inductance values are calculated, optimum number of turns and the coil dimensions can be obtained for the whole system design. Due to large coil dimensions and high number of turns in EV applications, 3D FEA analysis leads to long solution times for obtaining the inductance values. In order to calculate these values fast and accurately, an analytical approach must be developed.

The mutual inductance between two single and identical rectangular turns as shown in Fig. 2 can be calculated with the formula given in Eq. (3). First, the magnetic flux density

equation at the point  $P(x_1, y_1, z_1)$  formed by the current passing through the primary coil must be obtained. Considering a differential length of this current as a small cylinder of length  $dl_1$ , the magnetic flux density equation can be obtained as given in (5). If the inner radius of the straight wire is  $r_w$  and the side lengths of the rectangular coils are 2a and the 2b, the total magnetic flux ( $\varphi_{\text{Total}}$ ) in the shaded area can be calculated as in (6). Thus, the mutual inductance formula of two rectangular coil with  $N_1$  and  $N_2$  turns can be written as in (7).

$$L_{12} = \frac{\phi_{12}}{I_1}$$
 and  $\phi_{12} = \oint_{S_2} \overrightarrow{B_1 dS_2}$  (3)

$$\overrightarrow{B_1} = \frac{\mu_0}{4\pi} I_1 \oint_{C_i} \frac{\overrightarrow{\mathrm{d}I_1}}{r^2} x \hat{r}$$
(4)

$$\overline{B_{1}} = \frac{\mu_{0}I_{1}}{4\pi} \cdot \frac{(y_{1} - b)}{\left[(y_{1} - b)^{2} + z_{1}^{2}\right]} \cdot \left[\frac{x_{1} + a}{\sqrt{\left((x_{1} + a)^{2} + (y_{1} - b)^{2} + z_{1}^{2}\right)}} - \frac{-x_{1} + a}{\sqrt{\left((x_{1} - a)^{2} + (y_{1} - b)^{2} + z_{1}^{2}\right)}}\right] \cdot \hat{k}$$
(5)

$$\varphi_{\rm T} = \int_{-a}^{a} \int_{-b}^{b} \overline{B_1} dx_1 dy_1 \tag{6}$$

$$M = N_1 N_2 \frac{\varphi_{\text{Total}}}{I} \tag{7}$$

In the third section, detailed steps for calculations and long form of Eq. (7) are given. A similar approach is followed when calculating the self-inductance of the rectangular coil. In contrast to the mutual inductance calculation, the magnetic flux density equation generated by the current passing through the coil is obtained at a point which is in its inner surface and the self-inductance value can be obtained by integration along this surface. The magnetic flux density equation at this point is given in (8).

$$\overline{B}_{1} = \frac{\mu_{0}I_{1}}{2\pi} \cdot \left[ a \ln \left( \frac{\left(a + \sqrt{a^{2} + r_{w}^{2}}\right)\left(b - r_{w}\right)}{r_{w}\left(a + \sqrt{a^{2} + \left(b - r_{w}\right)^{2}}\right)} \right) + \sqrt{a^{2} + \left(b - r_{w}\right)^{2}} - \sqrt{a^{2} + r_{w}^{2}} - b + 2 \cdot r_{w} \right]$$
(8)

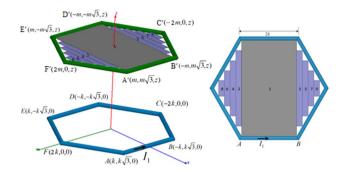


Fig. 3 a Placement of hexagonal coils on xyz coordinates. b Hexagonal coil divided into rectangular pieces

## 3 The proposed semi-analytical inductance calculation method for hexagonal coils

For the hexagonal coil structure, while calculating the magnetic flux to determine the self and mutual inductance, the coil inner area can be divided into rectangular and triangular pieces and the amount of flux in each part can be calculated separately. Then the total flux density can be obtained by summing them. However, to calculate the magnetic flux in the triangular pieces, some complex integration steps must be carried out. In this work, the triangular areas are divided

$$2k_{\rm avg} = \frac{2k_{\rm out} - [Nw + (N-1)g]}{2\sqrt{3}} \tag{9}$$

where w and g show the turn width and turn spacing, respectively.

First the main rectangular region along the length of the edge is created. Then, identical rectangular regions on the right and left sides of the main rectangular region are created. Thus, the regions numbered as 2-3, 4-5, 6-7 and 8-9 are identical. Therefore, it is sufficient to calculate the magnetic flux for areas 1, 2, 4, 6 and 8. As the number of rectangles increases, the amount of small triangular regions that are neglected decreases and the result obtained is more accurate. In order to find a more accurate inductance value, an iteration process has been applied to find the optimum number of rectangular regions in the triangle regions. For this, the height of the formed isosceles triangle is divided into identical parts and if their total number is n, a triangle is divided into rectangles with number of *i*. In this case, the sum of the calculated flux in the rectangular regions starting from *i* to n-i gives the total flux equivalent to the area of the triangular region. Equation (10) shows the flux calculation and integral boundaries for small rectangular areas obtained by dividing the triangular areas.

$$\varphi = \frac{\mu_0 I}{4\pi} \int_{-m\sqrt{3}\left(\frac{n-i}{n}\right)}^{m\sqrt{3}\left(\frac{n-i}{n}\right)} \int_{\left[\left(k\sqrt{3}-y_1\right)^2 + z_1^2\right]}^{m\left(\frac{n+i-1}{n}\right)} \cdot \left[\frac{k+x_1}{\left[\left(k\sqrt{3}-y_1\right)^2 + z_1^2\right]} + \frac{k-x_1}{\sqrt{\left((k-x_1)^2 + \left(k\sqrt{3}-y_1\right)^2 + z_1^2\right)}}\right] \cdot dx_1 \cdot dy_1$$
(10)

into small rectangular pieces as shown in Fig. 3 to calculate the inductance by using simple magnetic flux equations as described above. Then all the flux values are added up to find the total amount of flux. The magnetic flux density generated by the current flowing through the primary coil, in each defined area can be calculated as in (5).

2k and 2m show the side lengths of the single turn primary coil and secondary coil, respectively. For multi-turn spiral coils, the side length must be taken as the average value  $(2k_{avg})$  of the outer  $(2k_{out})$  and inner  $(2k_{in})$  side lengths for both coils by using simple trigonometric equations. For the primary coil, the calculation of the side length for multiturn spiral coil is, Figure 4 shows the flowchart describing the proposed algorithm to calculate the inductance values. In the flowchart, number of turns, side lengths, relative permeability of the air and radius of the conductor cross section area are defined as input values. In the first iteration, the flux is calculated for the values of n=3, i=1 and i=2 and the first inductance value is obtained by summing with the amount of flux in the large rectangular region-1. The second iteration runs for n=5 and i increments up to n-1 to obtain the new flux amount for triangular areas and by summing with the flux amount of region-1 a new inductance value can be calculated. The last inductance value is subtracted from the previous calculated inductance value to observe the percentage change. The iteration process continues the same way until the change between two consecutive steps is reduced to 0.001%. When the specified change is

reached, it is decided that no more division is necessary and the calculated inductance is the final value.

# 3.1 Mutual inductance

In order to calculate the mutual inductance between two hexagonal coils, the magnetic flux in the inner region of the secondary coil which is generated by the current flowing through the straight wire on one side length (AB) of the primary coil must be calculated. The flux amount in the inner area of secondary coil generated by the segment AB can be obtained by the sum of the fluxes of one rectangular and two triangle areas. In this case, the total flux ( $\varphi_{\text{total}}$ ) value will be six times the amount of flux obtained with a side length. The solution method for the calculation of the flux for two triangle regions  $\varphi_{{\rm AB},\Delta 1}$  and  $\varphi_{{\rm AB},\Delta 2}$  includes the same solution steps for the triangle areas which are given in "Appendix A" for the selfinductance calculation with the details. To calculate the flux amount in the area-1( $\varphi_{ABR1}$ ) which is shown in Fig. 3, the equation is divided in four parts (16)-(19) and the same solution steps are processed as described in Sect. 2.

$$\phi_{AB} = (\phi_{AB,\Delta_1} + \phi_{AB,\Delta_2} + \phi_{AB,R_1}) \tag{11}$$

$$\phi_{\text{Total}} = 6 \cdot \phi_{\text{AB}} \tag{12}$$

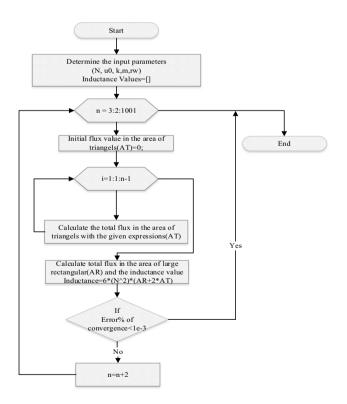


Fig. 4 Flowchart for inductance calculation

$$\phi_{AB,\Delta_1} = \frac{\mu_0}{2\pi} (S_1 - S_2 - S_3 + S_4)$$
(13)

$$\phi_{AB,\Delta_2} = \frac{\mu_0}{2\pi} \left( S'_1 - S'_2 - S'_3 + S'_4 \right) \tag{14}$$

$$\phi_{AB,R_1} = -\frac{\mu_0}{2\pi} \left( P_1 - P_2 - P_3 + P_4 \right) \tag{15}$$

$$P_{1} = \sqrt{\left(k\sqrt{3} - m\sqrt{3}\right)^{2} + z^{2} + (m+k)^{2}} - (m+k) \cdot \arctan h \frac{(m+k)}{\sqrt{\left(k\sqrt{3} - m\sqrt{3} + z^{2} + (m+k)^{2}\right)}}$$
(16)

$$P_{2} = \sqrt{\left(k\sqrt{3} - m\sqrt{3}\right)^{2} + z^{2} + (-m+k)^{2}} + (-m+k) \cdot \arctan h \frac{(-m+k)}{\sqrt{\left(k\sqrt{3} - m\sqrt{3} + z^{2} + (-m+k)^{2}\right)}}$$
(17)

$$P_{3} = \sqrt{\left(k\sqrt{3} + m\sqrt{3}\right)^{2} + z^{2} + (m+k)^{2}} + (m+k) \cdot \arctan h \frac{(m+k)}{\sqrt{\left(k\sqrt{3} - m\sqrt{3} + z^{2} + (m+k)^{2}\right)}}$$
(18)

$$P_{4} = \sqrt{\left(k\sqrt{3} + m\sqrt{3}\right)^{2} + z^{2} - (-m+k)^{2}} + (m+k) \cdot \arctan h \frac{(-m+k)}{\sqrt{\left(k\sqrt{3} - m\sqrt{3} + z^{2} + (-m+k)^{2}\right)}}$$
(19)

 $S_1$ - $S_4$ ,  $S'_1$ - $S'_4$  and  $P_1$ - $P_4$  show the part of the integrals in equations. As it can be seen from the derived equations, the mutual inductance formula is related to the side length of primary and secondary coils and the distance between these coils. Therefore, it is obvious the derived formula can be applied to coils with different dimensions and even structures.

#### 3.2 Self-inductance

Self-inductance calculations can be performed for the hexagonal coil structure with similar approach of the proposed method. First of all, the magnetic flux density formula is obtained for an arbitrary point in the inner area of hexagonal coil. By the integration of this formula through the surface enclosed by the side edges of the coil, the total magnetic flux can be determined. The magnetic flux density equation formed by the current flowing through the AB side length at any point in the divided rectangular parts as shown in Fig. 3b can be expressed in cylindrical coordinate system as given below.

$$\vec{B} = \frac{\mu_0 I}{4\pi} \frac{1}{r} \cdot \left[ \frac{z+k}{\sqrt{(z+k)^2 + r^2}} - \frac{z-k}{\sqrt{(z-k)^2 + r^2}} \right] \cdot \hat{\varphi} \quad (20)$$

As in the mutual inductance calculation, the total magnetic flux in the two identical triangles and a large rectangular area is calculated to obtain the self-inductance value. With the same approach, triangular areas are divided into rectangles to calculate the magnetic flux of the triangular areas with the least error rate and are added to magnetic flux within the large rectangular area to obtain the total magnetic flux. Derivation of the self-inductance formula is given in "Appendix A" in detail.

## 3.3 Verifications

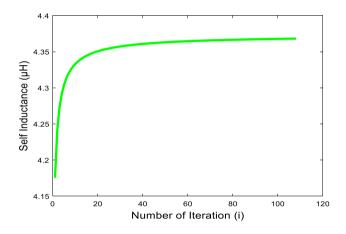
In this section, inductance calculation approaches presented in the literature are briefly explained and the results obtained by the proposed method are compared with their results. The semi-analytical approach proposed in this work can be used to determine both self and mutual inductances. Three simple approaches for calculating the self-inductance of square, circular, hexagonal and octagonal coils are given in Eqs. (21)–(23) [23, 27, 28]. Mohan et al. also modified the Wheeler's self-inductance formula and here  $\rho$ , N and  $d_{ave}$ show the fill ratio, number of turns and average diameter of the coil, respectively. In another simple approach, each side length of a polygonal coil is considered as a single current plate with the same current density. The common inductances of the current plates perpendicular to one another are zero. The self and mutual inductances are also determined by the geometric mean difference (GMD), the arithmetic mean difference (AMD) and the arithmetic mean squared difference (AMSD). The accuracy of the equation in (15) depends on the distance between the spiral coils (s) and the turn width (w). However, in practice, since the distance between the coils is very small, the equation works good. Equation (23) has been developed using the data matching technique and may not provide accurate results for new designs as it is resolved using the existing data library.

$$L_{\rm mw} = K_1 \mu_0 \frac{N^2 d_{\rm avg}}{1 + K_2 \rho}$$
(21)

$$L_{\rm gmd} = \frac{\mu N^2 d_{\rm avg} c_1}{2} \left( \ln(c_2/\rho) + c_3 \rho + c_4 \rho^2 \right)$$
(22)

Table 1 Comparing the proposed method with other analytical approaches

Method (outer side length)	Self-inductance (µH)		
	$(2k_{out}=10 \text{ cm})$	$(2k_{\text{out}}=20 \text{ cm})$	
Grover's Formula [11]	4.52	10.62	
Data Fitted Monomial Expression [23]	4.93	12.01	
Current Sheet Expression [28]	4.47	10.57	
Modified Wheeler Expression [27]	3.98	8.51	
Proposed Method	4.37	10.37	
Measurement Result	4.64	10.67	
Error (%)	5.8	2.8	



**Fig. 5** The variation of the self-inductance value with the number of iterations (i)

$$L_{\rm mon} = \beta d_{\rm out}^{\alpha_1} w^{\alpha_2} d_{\rm avg}^{\alpha_3} N^{\alpha_4} s^{\alpha_5}$$
(23)

The proposed semi-analytical self and common inductance calculation method was compared with other calculation methods used in the literature, for coils with 1.78 mm<sup>2</sup> litz wire cross section as three number of turns and for two different hexagonal coil with 10 and 20 cm outer side lengths. As can be seen in Table 1, the self-inductance values calculated with the proposed method are very close to the results obtained in the comparative studies. The percent error values in the tables from 1 to 4 show the results obtained by comparing the results of the proposed method with the experimental measurements.

In the calculation of the self-inductance of the hexagonal coil with an outer side length of 10 cm, the approach of the proposed method to the self-inductance value with the convergence error rate of 0.001% recommended for solution as a result of 108 iterations is shown in Fig. 5.

The analytical formulations proposed by Grover for the mutual inductance calculation between the identical polygonal coils for radio circuits are still one of the most frequently

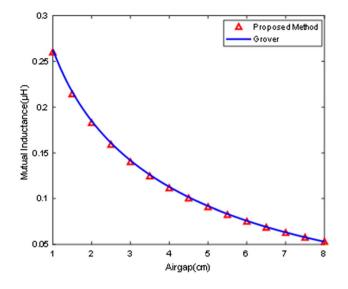


Fig. 6 Comparison of mutual inductance calculation

used formulations today. These formulas, developed using geometric and arithmetic mean differences, are getting more complex with using different coefficients in calculating the inductance of polygonal coils of different sizes and with large air gaps. In EV applications, the primary coil fixed to the platform on the ground generally does not have a size limitation, while the dimensions of the secondary coil fixed to the vehicle chassis are determined by the manufacturer and generally smaller than the primary coil [26]. In addition, different formulas need to be applied for different polygons in Grover's method and Eq. (24) shows the formula for the mutual inductance calculation of two identical hexagonal coils. Here *s* and *d*, respectively, represent the side length and the distance between the coils in centimeters [11].

$$M = 0.012s \left[ \log\left(\frac{s}{d}\right) - 0.151524 + \frac{d}{s} \left(1 - \frac{\pi\sqrt{3}}{9}\right) + \frac{d^2}{s^2} \left(\frac{2\sqrt{3} - 3}{4}\right) - \frac{d^4}{s^4} \left(\frac{33\sqrt{3} - 46}{216} + \cdots\right) \right]$$
(24)

Since the Grover method can be used to calculate the mutual inductance between identical coils, the proposed method has been applied for two identical and single turn hexagonal coils with distances ranging from 1 to 8 cm between the primary and secondary coils of 20 cm side lengths. As shown in Fig. 6, the results of the proposed method well match with the Grover's formula results. Grover's extended formula gives more accurate results than its simple form which is given in (24). In addition, the approximation to the mutual inductance between the identical hexagonal coils with three number of turns, 10 cm air gap and

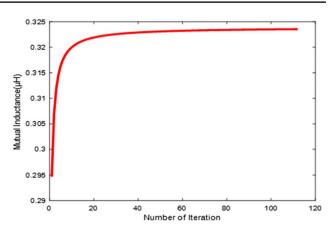


Fig. 7 The variation of the mutual inductance value with the number of iterations

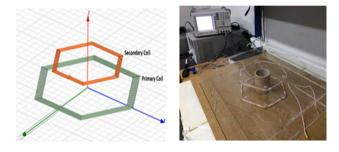


Fig. 8 3D FEA model and the manufactured hexagonal coils

 
 Table 2
 Comparison of FEA and experimental results with proposed method for self-inductance calculation

Self-inductance (µH)	$2k_{\text{out}} = 10 \text{ cm}$	$2k_{\text{out}}=20 \text{ cm}$
Measured	4.64	10.67
FEA	4.51	10.54
Proposed Method	4.37	10.37
Error (%)	5.8	2.8

10 cm outer side lengths is obtained with an acceptable error rate which is a result of 112 iterations and given in Fig. 7.

# 4 Comparison of experimental, FEA and analytical results for different coil dimensions

As seen in the third section, the results of the methods in the literature are similar with the proposed method. In this section, the results obtained with the proposed method were compared with FEA and experimental results for different coil dimensions. For this purpose, two hexagonal coils each with three turns in 10 and 20 cm outer side length dimensions are manufactured. The presentation of the coils manufactured in the laboratory and 3D model of the coils created with ANSYS Maxwell software are given in Fig. 8. Table 2 shows the comparison of the results obtained, with FEA and measured self-inductance values with the LCR meter in order to verify the self-inductance calculations of the manufactured coils. The frequency of the LCR meter was set to 20 kHz which is a common operating frequency for wireless power transfer systems. The LCR meter has 0.1% measurement accuracy as stated in the User Manuel. That means 0.01  $\mu$ H error margin for a measurement of a coil with 10  $\mu$ H value.

When the results obtained in Table 2 are examined, while the self-inductance values of the coil structures in the given dimensions are required to long calculation time by FEA, the recommended method can be calculated self-inductance values with small error rates as 2.8 and 5.8%, respectively. Considering the results obtained in this section and section three, it is thought that the proposed approach will have an important place in the literature because of the fact that the equations can be calculated with small error rates and the equations used have simple solutions.

In order to verify the mutual inductance calculations obtained by the proposed method, the distance between identical coils with three turns was changed from 2 to 12 cm, and the results were compared with the FEA and the measured results. Mutual inductance measurement procedure is as follows; First, the self-inductance values of the primary and secondary coils are measured with LCR meter and summed up. Then, one end of the primary coil and one end of the secondary coil are connected in series and a value is measured with the LCR meter. The summation of the selfinductances is subtracted from this and divided into two to get the mutual inductance value. In addition, the number of iterations required to obtain results at a predetermined error rate for each distance is given in Tables 3 and 4.

It can be seen obviously from Tables 3 and 4 that the proposed method results well match with the measurement results. In EV applications, the air gap between the coils is generally 10 cm and for this air gap the mutual inductance

Table 3 Mutual inductance calculation and comparison of results  $(2k_{out} = 10 \text{ cm})$ 

Air gap (cm)	<i>M</i> (μH)		Number of	Error (%)	
	Measured	FEA	Numerical	iterations	
2	1.55	1.53	1.58	94	1.9
4	0.94	0.93	0.96	100	2.08
6	0.62	0.61	0.64	104	3.22
8	0.42	0.41	0.44	108	4.76
10	0.3	0.29	0.32	112	6.67
12	0.21	0.2	0.23	114	9.5

Table 4 Mutual inductance calculation and comparison of results  $(2k_{out} = 20 \text{ cm})$ 

Airgap (cm)	Μ (μΗ)			Number of	Error (%)
	Measured	FEA	Numerical	iterations	
2	4.48	4.6	4.63	91	3.50
4	3.12	3.2	3.25	94	3.63
6	2.38	2.41	2.48	97	4.13
8	1.87	1.89	1.97	100	5.31
10	1.51	1.52	1.6	102	5.84
12	1.24	1.25	1.32	104	6.29

value for two identical hexagonal coils with 20 cm outer side lengths is calculated with just 5.84% error rate when compared to actual measurement results. When iteration number is examined, it is necessary to have higher iteration numbers in order to obtain the same error rate in small size coils. The reason for this is that the ratio of neglected area value to total area has a more significant value in small size coils. As the distance between the coils increases, the increase in the percentage error of the calculated mutual inductance value can be explained by the increasing of the amount of leakage at the same time. Figure 9 shows the comparison of the obtained mutual inductance results for the proposed method, FEA and experimental measurements at different air gaps with 20 cm side lengths of two hexagonal coils. When taking into consideration the measurement conditions, the number of turns which is more than one and the occurrence of parasitic errors due to the connectors, the inductance calculations can be performed by using the new proposed method with small error rates.

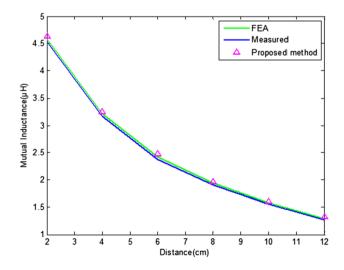


Fig. 9 Comparison of results

#### 5 Conclusion

In this study, a new method has been proposed by using simple magnetic flux distribution formulas which can be used for self and mutual inductance calculations of hexagonal coil structures. In order to prove the accuracy of the proposed method, the obtained results are compared with experimental measurements. FEA and different solution methods. The self-inductance value of the hexagonal coil with a side length of 20 cm calculated by the proposed method can be obtained with 2.8% error when compared to the experimental results. The distance between the coils is generally 10 cm in EV applications, and it is seen that the mutual inductance value can be obtained with 5.84% error for this distance when the proposed method result is compared with the experimental measurements. Considering that the inductance values are also large due to the size of the coils in these high-power applications, it is acceptable that these error rates will not cause a significant change in the number of turns. There are different approaches for inductance calculations in the literature, and some of them have great accuracy, especially for the multi-turn and small-sized coils. However, they are using complex integrals and different coefficients which is a disadvantage for these methods. Since with a simple approach, both the self and mutual inductance values can be obtained with acceptable error rates, the proposed method is a good candidate for the calculation of inductance values. For the future work, the proposed method can be used for inductance calculations of all polygons.

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# Appendix

Magnetic flux density equation can be written by the Biot-Savart Law,

$$\vec{B} = \frac{\mu_0 I}{4\pi} \frac{1}{r} \cdot \left[ \frac{z+k}{\sqrt{(z+k)^2 + r^2}} - \frac{z-k}{\sqrt{(z-k)^2 + r^2}} \right] \cdot \hat{\varphi} \quad (25)$$

The total magnetic flux equation for calculating the triangle areas which are divided into rectangles, depending on the side length of the hexagon and the division parameters n and i is can be obtained as follows,

$$\varphi_{AB} = \frac{\mu_0 I}{4\pi} \int_{\frac{k\sqrt{3}(2n-i)}{n}}^{\frac{k\sqrt{3}(2n-i)}{n}} \int_{\frac{k(n+i-1)}{n}}^{\frac{k(n+i-1)}{n}} \frac{z+k}{\sqrt{(z+k)^2 + r^2}} - \frac{z-k}{\sqrt{(z-k)^2 + r^2}} \cdot dz \cdot dr$$
(26)

ka/2(2m i)

$$\varphi_{AB} = \frac{\mu_0 I}{4\pi} \int_{\frac{k\sqrt{3}i}{n}}^{\frac{k\sqrt{3}i}{n}} \frac{1}{r} \cdot \left[ \sqrt{\left(\frac{k(2n+i)}{n}\right)^2 + r^2} - \sqrt{\left(\frac{ki}{n}\right)^2 + r^2} - \sqrt{\left(\frac{ki}{n}\right)^2 + r^2} - \sqrt{\left(\frac{k(2n+i-1)}{n}\right)^2 + r^2} + \sqrt{\left(\frac{k(i-1)}{n}\right)^2 + r^2} \right] \cdot dr$$
(27)

In order to simplify this complicated integration equation, we can divide it in four pieces and solve each part of it. As an example, only the solution of part  $A_1$  is given and the other parts can be solved by the same approach.

$$\varphi_{AB,\Delta} = \frac{\mu_0 I}{4\pi} \left( A_1 - A_2 - A_3 + A_4 \right)$$
(28)

$$A_{1} = \sqrt{\left(\frac{k(2n+i)}{n}\right)^{2} + \left(\frac{k\sqrt{3}(2n-i)}{n}\right)^{2}} - k\left(\frac{(2n+i)}{n}\right)\ln\left[\frac{\sqrt{\left(\frac{k(2n+i)}{n}\right)^{2} + \left(\frac{k\sqrt{3}(2n-i)}{n}\right)^{2}} + \left(\frac{k(2n+i)}{n}\right)}{\left(\frac{k(2n+i)}{n}\right)}\right]$$

$$(29)$$

Since there are two symmetrical triangles within the hexagonal area, the sum of the total flux formed by the AB length is equal to the sum of the two triangular areas. Thus, the total flux equation of the triangles is,

$$\varphi_{\mathrm{T,AB},\Delta} = 2\frac{\mu_0 I}{4\pi} (A_1 - A_2 - A_3 + A_4)$$
  
=  $\frac{\mu_0 I}{2\pi} (A_1 - A_2 - A_3 + A_4)$  (30)

Finally, the large rectangular area needs to be calculated. With the same approach, the total flux equation can be obtained if the integral boundaries are re-determined,

$$\varphi_{\text{AB,rectangular}} = \frac{\mu_0 I}{4\pi} \int_{r_w}^{2k\sqrt{3}-r_w} \frac{1}{r} \\ \cdot \left[ \int_{-k}^k \frac{z+k}{\sqrt{(z+k)^2+r^2}} \cdot dz - \int_{-k}^k \frac{z-k}{\sqrt{(z-k)^2+r^2}} \cdot dz \right] \cdot dr$$
(31)

After simplification of the equations for calculating the selfinductance of a hexagonal coil, the total flux is,

$$\varphi_{\rm AB,total} = 6 \left( \varphi_{\rm T,AB,\Delta} + \varphi_{\rm AB,rectangular} \right) \tag{32}$$

By using this equation, the self-inductance for a hexagonal coil for any size and wire radius is given by,

$$L_{\text{hexagon}} = \frac{6N^2 \left(2\varphi_{\text{AB},\Delta} + \varphi_{\text{AB,rectangular}}\right)}{I}$$
(33)

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